correspondence

Inelastic scattering puts in question recent claims of Anderson localization of light

To the Editor — In a recent Letter in Nature Photonics, Sperling et al.¹ reported the observation of Anderson localization of light in three dimensions. In contrast to previous studies, the authors designed their experiment to be insensitive to absorption. To this end, they measured the time-dependent broadening of a highintensity, short laser pulse transmitted through a highly scattering medium made from compressed TiO₂ powder. In analogy with the case of disordered optical fibres², localized light is expected to be laterally confined to roughly the localization length ξ . Using an ultrafast imaging system, Sperling et al. experimentally observed the saturation of the time-dependent transverse width of the total transmitted light intensity, and from this, they derived the claim for the first unequivocal observation of the three-dimensional localization of light.

In this correspondence, we would like to point out that the Letter of Sperling *et al.* does not report on the observation of elastic scattering of light waves, which is considered to be a necessary condition for the occurrence of Anderson localization. In his recent PhD thesis³, Wolfgang Bührer (who was supervised by Maret and Aegerter) reported that highly nonlinear contributions exceed the linear (elastic) scattering signal by at least one order of magnitude in the most relevant regime, namely the long-time regime ($\tau/\tau_{max} \approx 3$ in Fig. 2 of ref. 1). Although such

Maret et al. reply: The interplay between nonlinear effects and Anderson localization in disordered optical fibres¹ has recently attracted great interest, and it is important in the action of random lasers in which closed multiple scattering loops have enhanced intensity². As optical nonlinearities in TiO₂ can give valuable information on the nature of light transport in strongly scattering powders, we studied these effects in an extended experimental and theoretical investigation (to be published). Now, Scheffold and Wiersma have put forward that such effects may question the interpretation of the results of our recent experiments in terms of Anderson localization^{3,4}.

incoherent light makes an extremely small contribution to the total transmission, it becomes the dominant contribution to the late arriving signal once the elastically scattered light has leaked out of the sample. Wavelength-resolved experiments reported in this thesis on samples similar or identical to the ones studied in ref. 1, show that the non-exponential tail of the transmitted pulse disappears when the spectrally shifted contributions are blocked by a band-pass filter.

We believe that the dominant contribution of incoherent light puts in question not only the recent claims by Sperling *et al.* but also similar claims of localization by the same group in 2006 based on time-resolved measurements alone⁴. The long-time regime lies at the heart of both claims of localization, as this is the regime in which the saturation of the transverse width was observed in ref. 1 and in which the deviation from nonexponential decay was observed in ref. 4. Although Bührer extensively studied this incoherent contribution previously³, it was not mentioned in ref. 1.

We note that it is relatively easy to confuse nonlinear effects with localization in this type of experiment using pulsed laser sources. For example, photons generated by nonlinear processes (such as radiative decay after two-photon absorption) are emitted with a distribution of time delays Δt , which contribute to a narrower transmission profile

As noted by Scheffold and Wiersma, our measurements of the time-dependent transmission profiles4 eliminate the influence of absorption. This technique thus reveals the signatures of Anderson localization more clearly than integrated transmission data³. The relevant signals appear at different timescales, which is also true for the nonlinear contributions in our study to be published. Based on the arguments put forth by Scheffold and Wiersma, one would expect that no saturation of the profile would be observed when a band-pass filter for the incoming wavelength is inserted. However, such an experiment (Fig. 1) clearly shows

 $T(\rho, t) = T_{\text{elastic}}(\rho, t) + T_{\text{inelastic}}(\rho, t - \Delta t)$ at the output. The transition from elastic to inelastic scattering in the long-time regime can thus result in an 'apparent' saturation of the transverse width, resembling that of localization. Also, the different particle sizes and the different pressures used to produce samples of varying packing fractions can lead to differences in nonlinear optical coefficients, which can be easily misinterpreted as a localization effect that depends on the scattering strength (*kl**), whereas it is actually a nonlinear optical effect.

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saturation of the profile width, similar to that reported in ref. 4, irrespective of the detected frequency.

In addition, we emphasize that the dependence of the localization length on kl^* was determined in two separate ways in ref. 4. Besides the variation of the samples and the packing fractions mentioned by Scheffold and Wiersma, we also investigated the change in the scattering strength induced by varying the incident wavelength. Our data presented in two different studies^{3,4}, exhibit a remarkable scaling with the turbidity measured by kl^* — exactIZ as expected for three-dimensional Anderson localization.

Even if, in an extraordinary coincidence, the change in nonlinear effects mentioned by Scheffold and Wiersma scaled exactly as kl^* , an alternative method of changing kl^* would not give the same quantitative results, whereas we have observed that it does⁴. Also, the determination of the localization length in two experiments that are conceptually very different^{3,4} gives the same results (within errors) as a function of kl^* .

Although in different experiments^{3,4}, the input power density varied by more than four orders of magnitude, the features in the time-dependent integrated transmission varied by less than a factor of five at long times, showing that the nonlinear contributions are weak. Moreover, the nonlinear scenario invoked by Scheffold and Wiersma to explain the narrowing of the observed profiles would not cause saturation at long times, rather they would only reduce the rate of increase, which has not been observed. In that case, the relative contributions should be strongly dependent on the incident intensity, which again has not been observed.

In conclusion, elastically scattered light does show a plateauing width in our experiments, and the observation of a plateau clearly scales with kl*, as determined by two independent means. Small nonlinear contributions via phononassisted processes do exist and provide density of states inside the electronic bandgap of TiO₂, which can qualitatively account for the observed frequency shifts given the intensity enhancement in the localized modes. Both integrated timedependent transmission and transmission profiles are quantitatively accounted for by the self-consistent Vollhardt-Wölfle theory of Anderson localization of photons in the presence of realistic nonlinear effects,

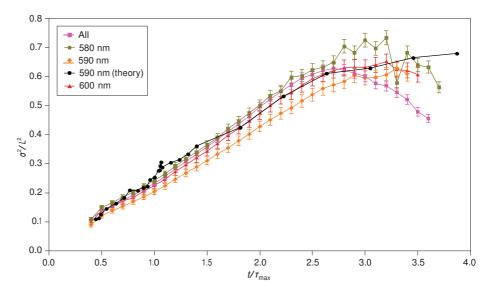


Figure 1 | Time dependence of the mean square width σ^2 normalized to sample thickness *L* for pressed powder of a strongly scattering sample (R700, see ref. 4, $kl^* = 2.7(3)$, L = 0.86 mm, absorption time = 0.9(1) ns) at incident laser wavelength 590 nm and different band-pass filters between the sample and the detector. The legend shows the centre filter wavelengths; the filter width was 10 nm. Coloured lines represent experimentally measured data, whereas the black line indicates curve predicted by theory for the same sample parameters as those used in the experiment.

boundaries and other sample parameters⁵ matched to the experiment (see Fig. 1). Thus, although small nonlinear effects may eventually be present, they do not invalidate the experimental observation of the Anderson localization transition with light.

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